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To cite this article: Anton V Ievlev et al 2018 Nanotechnology 29 155302

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Received 12 December 2017, revised 15 January 2018
Accepted for publication 2 February 2018
Published 22 February 2018

Abstract

The ability to control thin-film growth has led to advances in our understanding of fundamental physics as well as to the emergence of novel technologies. However, common thin-film growth techniques introduce a number of limitations related to the concentration of defects on film interfaces and surfaces that limit the scope of systems that can be produced and studied experimentally. Here, we developed an ion-beam based subtractive fabrication process that enables creation and modification of thin films with pre-defined thicknesses. To accomplish this we transformed a multimodal imaging platform that combines time-of-flight secondary ion mass spectrometry with atomic force microscopy to a unique fabrication tool that allows for precise sputtering of the nanometer-thin layers of material. To demonstrate fabrication of thin-films with \textit{in situ} feedback and control on film thickness and functionality we systematically studied thickness dependence of ferroelectric switching of lead-zirconate-titanate, within a single epitaxial film. Our results demonstrate that through a subtractive film fabrication process we can control the piezoelectric response as a function of film thickness as well as improve on the overall piezoelectric response versus an untreated film.

Keywords: ferroelectrics, thin films, ion beam fabrication, time-of-flight secondary ion mass spectrometry, atomic force microscopy

(Some figures may appear in colour only in the online journal)

Introduction

The ability to control thin-film growth has led to advances in our understanding of fundamental physics as well as to the emergence of novel technologies including giant magnetoresistance materials and devices \cite{1–5}, ferroelectric tunneling barriers and capacitors \cite{6–9}, superconductor oxides of Josephson junctions \cite{10, 11}, quantum computing \cite{12–14}, and more exotic systems. The rapid improvement of modern electronics hinges on the development of novel methodologies to process materials at ever shrinking dimensions. Processing techniques that allow for monolayer precision control during the growth process like pulsed-laser deposition \cite{15} and atomic-layer deposition \cite{16} grow materials in accordance with fundamental physics and often provide only certain terminations at the film interface and surface.
Additionally, when using these approaches dislocations and vacancies will migrate to the interface of the films when high enough temperatures are reached causing problems for controlling the chemistry of the interface as well as leading to surface instabilities and contamination [17]. Altogether, these factors may hamper fundamental investigations and practical applications of thin film systems we can produce.

Here we open a new direction for research in complex oxides by demonstrating a subtractive fabrication process that enables creation and modification of thin films with pre-defined thicknesses. It allows us to elucidate the properties of oxides while simultaneously minimizing the effects of contaminants. We apply a combination of time-of-flight secondary ion mass spectrometry (ToF-SIMS) [18–21] with atomic force microscopy (AFM) for precise sputtering of the nanometer-thin layers of material. The sputtering process is realized by the ion beam under control of mass spectrometry. AFM is further utilized for sputtering depth calibration. To verify this technique, we systematically investigated thickness dependence of ferroelectric switching of lead-zirconate-titanate, with just one epitaxial film as a source of materials. Surprisingly, we observe not only robust switching in pre-sputtered films, but also strong improvement of piezoelectric response once the aged surface layer is removed.

**Experiment and results**

As a model system, we used a 100 nm thick ferroelectric film of lead zirconate titanate PbZr_{0.2}Ti_{0.8}O_{3} (PZT) grown on a SrTiO_{3} (STO) (001) substrate with a 50 nm thick SrRuO_{3} (SRO) buffer layer. ToF-SIMS was used to clean the PZT surface and adjust its thickness. Measurements have been performed using a TOF.SIMS.5 NSC (ION-TOF GmbH) instrument, combining ToF-SIMS and AFM in the same vacuum chamber. Mass spectrometry measurements have been done by a focused O_{2}– sputtering source (500 eV energy, 100 nA current, and ∼20 μm spot size) and a Bi+ primary gun (30 keV energy, 30 nA current, and ∼5 μm spot size). Sputtered regions were further analyzed in a built-in AFM chamber to measure the absolute values of the sputtering depth. External AFM (Icon Bruker) and STEM (Titan, FEI) were also used to compare crystallography and ferroelectric properties of pristine and treated regions.

As a first step, we performed depth profiling of the PZT locally, which revealed the total film thickness, and the local chemical composition of surface layer and the film. We used O_{2}– sputtering over 300 × 300 μm regions with cycle of 10 s. Each cycle was accompanied by ToF-SIMS surface chemistry analysis using the focused Bi+ primary ion source and the time-of-flight analyzer in positive ion-detection mode. The mass spectra are therefore collected from each layer, enabling depth profiling of the chemical composition.

The averaged mass spectrum of the PZT film and part of SRO buffer (figure 1) clearly shows peaks of the PZT base elements (Pb+, Zr2+ and Ti3+), their oxides (PbO+, ZrO2+ and TiO2+), Sr2+ and Ru3+ from the SRO buffer layer as well as elements from an adsorption surface layer (Na+, K+, etc). The area of those peaks has been used to plot depth profiles of their bulk distribution (figure 2(a)). One should note that the x-axis in figure 1 is not the physical depth, but the sputter time, which is directly proportional to the sputter depth (the longer sputtering time the deeper the crater). These profiles show sharp transition from PZT to SRO (decrease of Pb+, Zr2+, Ti3+ and increase of Sr2+ signals) at around 840 s of sputtering. This can be used to roughly estimate the sputtering rate of PZT: \(v_{\text{PZT}} \sim 0.12 \pm 0.01 \text{ nm s}^{-1}\).

We further performed precise calibration of the sputtering depth using the AFM built-in to the ToF-SIMS vacuum chamber. The AFM measurements were performed in contact mode inside the regions sputtered for different times from 210 to 930 s. The crater topography profiles were used for sputtering depth calibration (figure 2(b)). These measurements were carried out in automated profiling mode. In this mode resulted crater profile (300 μm long) is measured by sections of 80 μm. Sections are measured with overlap of 15% and further automatically stitched. Experiments demonstrated that the sputtering is linear with time within the film, and the value of the sputtering rate based on the AFM data is identical to that estimated from the ToF-SIMS depth profiles. In addition, we used macroscopic data of the film thickness and measured profile for the longest sputter to estimate the sputtering rate of the SRO buffer layer. The obtained value \(v_{\text{SRO}} \sim 0.028 \text{ nm s}^{-1}\) was found to be significantly lower than the sputtering rate of PZT. Hereinafter we will reference sputtering depth in percentage from initial film thickness, based on the sputtering time and the macroscopic data of the film thickness (figure 2(a)). Analysis of the depth profiles also showed slight increase in the roughness of sputtered region (figure 2(d)). Root mean squared \(R_{q}\) value of the roughness is used. However, total roughness change does not exceed 1.5 nm, which cannot significantly affect properties of the modified sample.

The results of this work reveal that the sputtering process by ToF-SIMS under the control of mass spectrometry can act as a subtractive fabrication approach for the creation of thin films of pre-defined thickness and surface conditioning. We should note, that works on the ion beam thinning of ferroelectric thin films and single crystals were performed in the past, however those approaches led to damage of the crystallographic structure and degradation of ferroelectric properties [22–24]. To ascertain structural damages during this process and establish potential ionic damage we performed STEM and AFM measurements.

STEM was used to investigate the atomic structure of cross-sections of pristine and sputtered PZT samples (figure 3). In all cases these measurements clearly differentiate the surface layer, the PZT film, and the SRO/STO substrate. The pristine, as-grown sample, as expected, showed the presence of a thick (∼10 nm) inhomogeneous aged surface layer (figure 3(a)). This layer is known to hamper local probing of ferroelectric properties by AFM and can lead to a number of non-trivial phenomena during polarization reversal [25–29]. Presence of this layer is confirmed by ToF-SIMS, which showed high concentration Na+, K+ on the sample surface (figure 2(a)). The sputtering process removed this
contamination layer (figures 3(b), (c)) and demonstrated consistency with designated sputtering times—the film thickness was ∼52 nm after sputtering for 50% (figure 3(b)), and ∼15 nm for 80% (figure 3(c)). Detailed STEM imaging showed that sputtering process does not change crystallography in the bulk of PZT and on the PZT/SRO interface (figure 3(d)), however it showed formation of the thin (∼3 nm) amorphous layer (figure 3(e)). At this point, we cannot speculate about the origin of this amorphous layer, it could be formed during the sputtering process or just be an artifact of the cross-section preparation for STEM imaging.

We should note here, that thickness of this artificial layer is much smaller, than in the case of pristine PZT sample, so one would expect enhancement in the properties of modified films.

We further used the AFM to study ferroelectric and conductive properties of the treated regions and piezoresponse force microscopy was used to study switching. The measurements were carried out over a grid of 10 × 10 points with an area of 5 × 5 μm.

We used local hysteresis loops to characterize the ferroelectric properties of the sputtered PZT regions (figure 4).
These loops are measured during application of the DC voltage to the tip and are able to characterize ferroelectric switching in the nanometer region underneath the tip under the action of external DC field produced by the tip. Details about measurements of the local hysteresis loops and their interpretation can be found in [30–35].

Careful analysis of the averaged loops measured at the sputtered regions clearly showed hysteresis, which confirms that the film remains ferroelectric, as polarization can be switched by application of an external electric filed. Furthermore, the coercive voltage (loop width on the half max) was not significantly altered, ranging from 4 to 5 V for all regions. Coercive voltage corresponds to the electric field required to reverse polarization underneath the tip and is one of the most important ferroelectric parameters.

At the same time, we found, that the total piezoresponse (full height of the loop), has been significantly affected. It increased more than twice from 31 mV for the pristine, as-grown sample to 71 ± 8 mV after a short sputtering process (1% of the film thickness) and to 90 ± 10 mV after sputtering of half the film thickness. We should note here that in this manuscript we measure piezoresponse in volts (mV), which is a deflection voltage reading from the AFM detector, which is directly proportional to deflection of the tip and piezoelectric deformation of the surface.

Increase in the total piezoresponse is in a good agreement with the STEM results, showing the presence of a significant adsorption layer on the surface of the pristine, as-grown sample (figure 3(a)). In this case, piezoresponse of the pristine sample is reduced due to the presence of the adsorption layer. This leads to decrease in the sensitivity of PFM measurements and correspondingly decreases piezoresponse. At the same time, ferroelectric properties of the film stay unchanged, which is confirmed by the loop shape and values of coercive field.

To characterize the conductive properties of the modified regions we performed current measurements through the AFM tip (figure 5). These measurements did not show any significant changes in the film electro-conductivity up to 90% of the sputtering (figure 5(a)). The corresponding resistivity was around 1–10 GΩ. However, further decreasing the residual film thickness down to 10% of the initial thickness led to...
abrupt increase of the current (to 100 nA) with corresponding resistivity $\sim 1 \, \text{M}\Omega$ (figure 5(b)). This phenomenon is related with the proximity of the conductive SRO buffer layer and demonstrates the challenge of obtaining extremely thin (few unit-cell) films with the sputtering approach.

**Methods**

**Time-of-flight secondary ion mass spectrometry**

ToF-SIMS measurements have been done using TOF-SIMS.5-NSC instrument, using bismuth ion gun as a primary ion source and $\text{O}_2^-$ ion gun as a sputter source. Sputter was realized by $\text{O}_2^-$ ion beam with energy of ions 1 keV, current 120 nA and spot size $\sim 20 \, \mu\text{m}$. To control sputtering process we analyzed mass spectra of the freshly sputtered surface in positive SIMS mode. In this case Bi liquid metal ion gun with energy of ions 30 keV, current $\sim 30 \, \text{nA}$ and spot size $\sim 5 \, \mu\text{m}$. This allowed us to acquire mass spectra with spectral resolution below 0.01 Da.

**Atomic force microscopy**

Characterization of the depth and profiles of sputtered craters has been done using atomic force microscopy Nanoscan introduced into the vacuum chamber of ToF-SIMS. Measurements were performed in contact AFM mode using Nanosensor PPP-EFM tips with force constant of 2.8 N m$^{-1}$ and radius of curvature better than 25 nm. Piezoresponse Force Microscopy and conductive AFM measurements were carried out in environmental conditions using Bruker Icon AFM and Budget Sensors 75E-G tips with force constant of 3 N m$^{-1}$ and radius of curvature below 25 nm. In PFM measurements AC voltage with amplitude of 0.5 V and frequency about 300 kHz was applied to the tip in addition to switching DC voltage $-5$ to $+5 \, \text{V}$. For local current measurements DC voltage was applied to the tip, while current has been measured using external current amplifier connected to the bottom electrode of the sample.

**Scanning transmission electron microscopy**

A Hitachi NB5000 FIB/SEM was used to prepare electron transparent foils of the PZT/SrTiO3 thin film sample for STEM imaging. To protect the PZT surface from ion beam damage during preparation, the sample was coated with $\sim 200 \, \text{nm}$ of carbon. Thinning to electron transparency begun with deposition of a 25 $\mu\text{m}$ long, 4 $\mu\text{m}$ wide, $\sim 500 \, \text{nm}$ thick tungsten capping layer to reduce 'curtaining' during the final FIB thinning of the specimen. A 40 kV focused beam of Ga ions, with a current of 19.5 nA, was used to mill material away down to the substrate to a depth of 4–5 $\mu\text{m}$. The sample was lifted from the bulk and transferred to a Cu Omniprobe half grid. A series of milling steps was then used to reduce the thickness of the lamella until it was electron transparent ($< 100 \, \text{nm}$). The thinning was started with a 20 kV, 0.11 nA beam and thinned the sample to 1 $\mu\text{m}$ thick. For the final milling step the beam parameters were changed to 20 kV, 0.11 nA and the sample was thinned to $< 100 \, \text{nm}$ thick. To minimize Ga implantation effects and amorphous material on the surface of the FIB thinned sample, a Fischione Nanomill was operated at 900 eV with a 130 pA Ar$^+$ beam for ion milling of the sample at $\pm 10^\circ$ for 15 min each side.

High angle annular dark field (HAADF) imaging of the PZT films was performed at 100 kV with an aberration corrected scanning transmission electron microscope, the Nion UltraSTEM 100. For HAADF imaging an illumination half angle of $\sim 85 \, \text{mrad}$ was used.

**Conclusion**

In conclusion, in this paper we developed a universal technique for subtractive fabrication of thin films with precisely controlled thickness. In this approach, a focused oxygen-ion beam is used for sputtering under the control of ToF-SIMS, which allows tuning of functional material response through in situ surface cleaning and feedback on precise film thickness. This approach can be useful for fundamental investigations as well as practical application of various thin
films. In particular, it enables size-effects on collective phenomena and correlated atomic and electron states to be investigated within a single film, modified in certain regions by the ion beam. Our benchmark study of ferroelectric size effect in lead zirconate-titanate revealed the presence of a significant adsorption layer on the surface of pristine, as-grown samples leading to a decrease of the piezoelectric response, measured with an AFM tip. Oxygen sputtering grown samples leading to a decrease of the piezoelectric effect in lead zirconate-titanate revealed the presence of a phenomena and correlated atomic and electron states to be

Acknowledgments

This material is based upon work supported by the U S Department of Energy, Office of Science, Office of Basic Energy Sciences under contract number DE-AC05-000OR22725 by the Laboratory Directed Research and Development Program of Oak Ridge National Laboratory, managed by UT-Battelle, LLC, for the US Department of Energy (AVL OSO). ToF SIMS, and AFM measurements were performed at the Center for Nanophase Materials Sciences, which is a DOE Office of Science User Facility. JCA acknowledges support from the Army Research Office under grant W911NF-14-1-0104. GAV acknowledges support from the National Science Foundation under grant DMR-1708615. LWM acknowledges support from the National Science Foundation under grant DMR-1608938.

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